Tomographic cross-correlation of the CMB lensing and galaxy clustering – systematic errors from cross-talk between redshift bins of galaxies

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CMB lensing

• Deflection of the CMB photon paths by the large scale structure of the Universe (${\sim}3')$

0.0016

• Correlation of deflection angles over the sky

-0.0016

- Reconstruction of lensing potential from perturbations of statistical properties of CMB anisotropy
- Tracer of dark matter distribution in broad redshift range

$$\phi(\hat{n}) = -\frac{2}{c^2} \int_0^{\chi_{rec}} d\chi \frac{D_{ls}}{D_l D_s} \Psi(\chi_0 - \chi, \chi \hat{n})$$





Planck collaboration et al. (2020)

Cross-correlation between CMB lensing and galaxy surveys

- Broad CMB lensing kernel does not allow tracing time evolution of dark matter
- Needed cross-correlation of CMB lensing map with objects with known redshift (galaxies, quasars, radio sources, etc.)
- Splitting redshift distribution on redshift bins (cosmic tomography: White et al. 2022; Pandey et al. 2022; Chang et al. 2022; Sun et al. 2022; Krolewski et al. 2021; Hang et al. 2021; Peacock & Bilicki 2018)





• Power spectrum of cross-correlation between CMB lensing and galaxy density contrast

$$\begin{split} C_{\ell}^{\kappa g} &= \int_{0}^{\chi_{*}} \mathrm{d}\chi \frac{W^{\kappa}(\chi) W^{g}(\chi)}{\chi^{2}} P_{m} \left(k = \frac{\ell + 1/2}{\chi}, z(\chi)\right) \qquad \theta \sim \frac{\pi}{\ell} \\ \kappa(\hat{\mathbf{n}}) &= -\frac{1}{2} \nabla^{2} \phi(\hat{\mathbf{n}}) \\ g &= \frac{n - \bar{n}}{\bar{n}} \\ W^{\kappa}(\chi) &= \frac{3\Omega_{m}}{2c^{2}} H_{0}^{2}(1 + z) \chi \frac{\chi_{*} - \chi}{\chi_{*}} \\ W^{g}(\chi) &= b(z(\chi)) \frac{H(\chi)}{c} \frac{\mathrm{d}N}{\mathrm{d}z(\chi)} \end{split}$$

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Ζ



- How well can we measure cross-correlations in different redshift bins given photometric redshift errors?
- Test using simulations of forthcoming galaxy photometric surveys (LSST)
- 300 simulations of correlated log-normal galaxy over-density (with LSST Science Book redshift distribution) and CMB lensing convergence fields (consistent with Planck CMB lensing map) using Full-sky Lognormal Astro-fields Simulation Kit (FLASK) code (Xavier et al. 2016)





Tomographic binning of redshift distribution

- Photometric redshifts $z_{\rm p}$ obtained by adding Gaussian photo-z errors to true redshifts $z_{\rm t}$

$$\frac{\mathrm{d}N(z_p)}{\mathrm{d}z_p} = \int \mathrm{d}z_t \frac{\mathrm{d}N(z_t)}{\mathrm{d}z_t} p(z_p - z_t | z_t) \qquad p(z_p - z_t | z_t) = G(z_t, \sigma(z_t)) \qquad \sigma(z) = \sigma_0(1+z)$$

• Tomographic binning of the true redshift distribution



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NAROBOWE CENTRUM BADAN JADROWYCH AWIERK

Power spectra for tomographic analysis

• Power spectra for galaxies with photometric redshifts are related to power spectra for galaxies with true redshifts by (Zhang et al. 2010):

$$C_{ij}^{gg,Ph}(\ell) = \sum_{k} P_{ki} P_{kj} C_{kk}^{gg,Tr}(\ell)$$

$$C_{i}^{\kappa g,Ph}(\ell) = \sum_{k} P_{ki} C_{kk}^{\kappa g,Tr}(\ell)$$
where $P_{ij} \equiv \frac{N_{i \to j}}{N_{j}^{Ph}}$ is so called scattering matrix

• Zhang et al. (2017) proposed algorithm, Non-negative Matrix Factorization, to solve

for P_{ij} and $C_{kk}^{Tr}(\ell)$ having $C_{ij}^{Ph}(\ell)$

• With estimation of the true redshift distribution it is possible fast method of computation of the scattering matrix $\sigma_0 = 0.02$ $\sigma_0 = 0.05$





• Naive approach to model power spectra for galaxies with photo-z

$$C_i^{gg,Ph}(\ell) = \int_0^{\chi_*} \frac{\mathrm{d}\chi}{\chi^2} \left(b(z_p) \frac{\mathrm{d}N^i(z_p)}{\mathrm{d}z_p} \right)^2 P_m\left(k = \frac{\ell + 1/2}{\chi}, z_p(\chi)\right)$$
$$C_i^{\kappa g,Ph}(\ell) = \int_0^{\chi_*} \frac{\mathrm{d}\chi}{\chi^2} W^{\kappa}(\chi) b(z_p) \frac{\mathrm{d}N^i(z_p)}{\mathrm{d}z_p} P_m\left(k = \frac{\ell + 1/2}{\chi}, z_p(\chi)\right)$$

• Estimation of the angular power spectra



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 $\sigma_0 = 0.02$



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 $\sigma_0 = 0.05$



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 $\sigma_0 = 0.05$



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Estimation of the parameters

Fitting parameters using the maximum likelihood method:

$$\mathcal{L}(\hat{C}_L|b,A) = \frac{1}{\sqrt{(2\pi)^{N_L} det(\operatorname{Cov}_{LL'})}} \times \exp\left\{-\frac{1}{2}[\hat{C}_L - C_L(b,A)](\operatorname{Cov}_{LL'})^{-1}[\hat{C}_{L'} - C_{L'}(b,A)]\right\}$$

• Linear galaxy bias: $b(z) = 1 + \frac{b_0 - 1}{D(z)}$ $C_L^{gg} \propto b(\bar{z})^2 C_L^{gg}$

• Amplitude of cross-correlation A: scales the amplitude of the cross-power spectrum (equals 1 for the ACDM model) $C_L^{\kappa g} \propto A \, b(\bar{z}) \, C_L^{\kappa g, {\rm fid}}$



Bin 5 $(0.8 \le z < 1.0)$

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 $\sigma_0 = 0.02$

Estimation of the parameters

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CosmoVerse@Lisbon, Jun 2023

 $\sigma_0 = 0.05$



Estimation of the parameters



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Estimation of the parameters



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Conclusions

- Tomographic cross-correlation between CMB lensing map and galaxy surveys useful for tracing time evolution of the large-scale structure
- Systematic errors caused by redshift bin mismatch of galaxies with photo-z
 - ~ ~10-15% for galaxy auto-power spectra, ~2% for cross-power spectra ($\sigma_0=0.02$)
 - ~1 standard deviation for the correlation amplitude, ~10-20 standard deviations for the galaxy bias ($\sigma_0 = 0.02$)
- Needed correction for the redshift bin mismatch using scattering matrix formalism
- Possible fast computation of the scattering matrix using estimation of the true redshift distribution (needed accurate estimation of the photo-z error distribution)